

Quantum Black Holes

theory & phenomenology

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II FLAG MEETING
Trento - June 7th, 2016



Netherlands Organisation for Scientific Research

Radboud University Nijmegen



(1971).
⁶ Bassi, P., Clark, G., and Rossi, B., *Phys. Rev.*, **92**, 441 (1953).
⁷ David, F. N., *Biometrika*, **34**, 299 (1947).
⁸ Weekes, T. G., *Nature phys. Sci.*, **223**, 129 (1971).

Black hole explosions?

QUANTUM gravitational effects are usually ignored in calculations of the formation and evolution of black holes. The justification for this is that the radius of curvature of space-time outside the event horizon is very large compared to the Planck length $(G\hbar/c^3)^{1/2} \approx 10^{-33}$ cm, the length scale on which quantum fluctuations of the metric are expected to be of order unity. This means that the energy density of particles created by the gravitational field is small compared to the space-time curvature. Even though quantum effects may be small locally, they may still, however, add up to produce a significant effect over the lifetime of the Universe $\approx 10^{17}$ s which is very long compared to the Planck time $\approx 10^{-43}$ s.

nations of the f_i and the \bar{f}_i :

$$p_i = \sum_j \{ \alpha_{ij} f_j + \beta_{ij} \bar{f}_j \} \quad \text{and so on}$$

The β_{ij} will not be zero because the time dependence of the metric during the collapse will cause a certain amount of mixing of positive and negative frequencies. Equating the two expressions for ϕ , one finds that the b_i , which are the annihilation operators for outgoing scalar particles, can be expressed as a linear combination of the ingoing annihilation and creation operators a_i and a_i^+

$$b_i = \sum_j \{ \bar{\alpha}_{ij} a_j - \bar{\beta}_{ij} a_j^+ \}$$

Thus when there are no incoming particles the expectation value of the number operator $b_i^+ b_i$ of the i th outgoing state is

$$\langle 0_- | b_i^+ b_i | 0_- \rangle = \sum_j |\beta_{ij}|^2$$

The number of particles created and emitted to infinity in a gravitational collapse can therefore be calculated by summing over all i and j .

For $m \sim 10^{24}$ Kg, $m^3 \sim 10^{50}$ Hubble times, while $m^2 \sim$ Hubble time

Nature

the collapse is spherically symmetric. The angular dependence of the solution of the wave equation can then be expressed in terms of the spherical harmonics Y_{lm} and the dependence on retarded or advanced time u, v can be taken to have the form $\omega^{-1/2} \exp(i\omega u)$ (here the continuum normalisation is used). Outgoing solutions $p_{lm\omega}$ will now be expressed as an integral over incoming fields with the same l and m :

$$p_\omega = \int \{ \alpha_{\omega\omega'} f_{\omega'} + \beta_{\omega\omega'} \bar{f}_{\omega'} \} d\omega'$$

(The lm suffixes have been dropped.) To calculate $\alpha_{\omega\omega'}$ and $\beta_{\omega\omega'}$ consider a wave which has a positive frequency ω on I^+ propagating backwards through spacetime with nothing crossing the event horizon. Part of this wave will be scattered by the curvature of the static Schwarzschild solution outside the black hole and will end up on I^- with the same frequency ω .

Beckenstein⁶ suggested on thermodynamic grounds that some multiple of κ should be regarded as the temperature of a black hole. He did not, however, suggest that a black hole could emit particles as well as absorb them. For this reason Bardeen, Carter and I considered that the thermodynamical similarity between κ and temperature was only an analogy. The present result seems to indicate, however, that there may be more to it than this. Of course this calculation ignores the back reaction of the particles on the metric, and quantum fluctuations on the metric. These might alter the picture.

Further details of this work will be published elsewhere. The author is very grateful to G. W. Gibbons for discussions and help.

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LOOP QUANTUM COSMOLOGY INSPIRATION

- Effective repulsive force

- Planck density

- Size \gg Planck length

- $r_b \sim \frac{m}{m_P} \ell_P$

- Quantum Tunneling

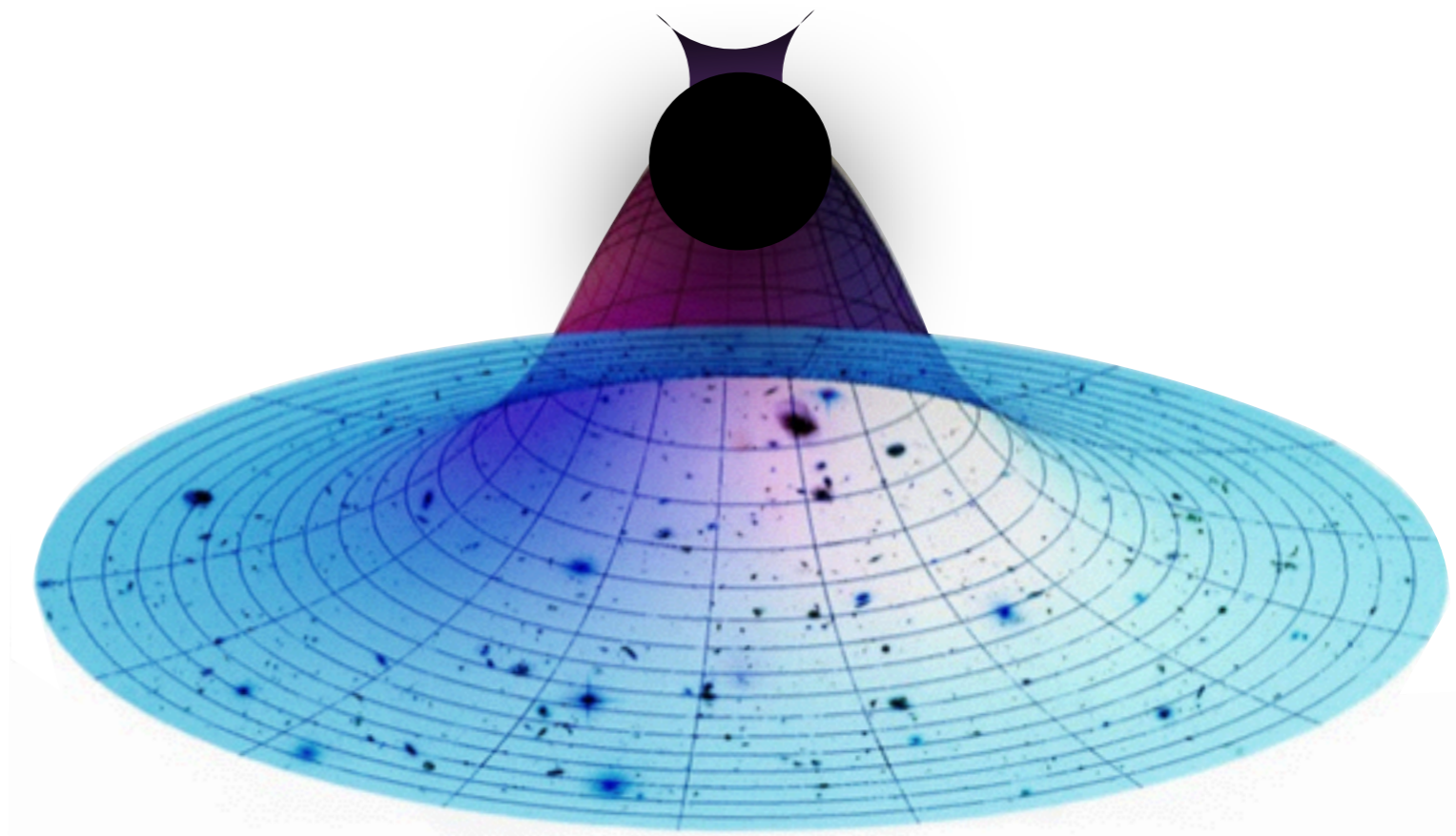
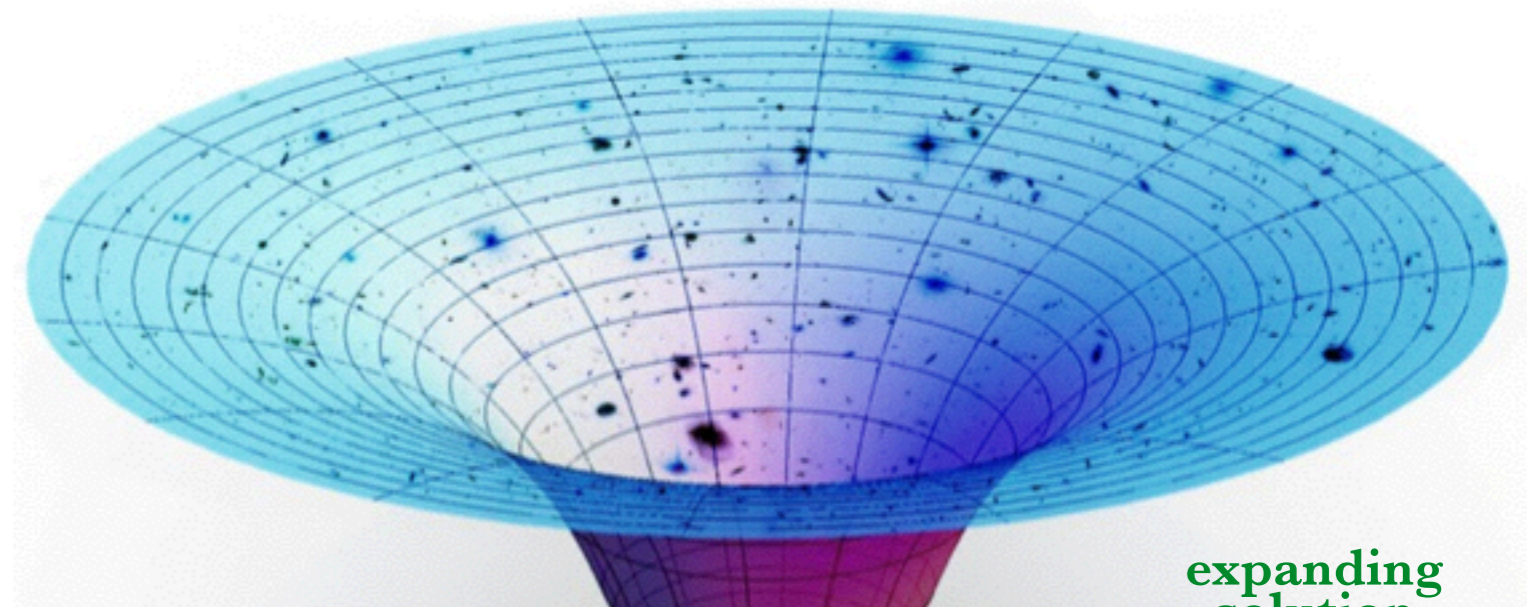
- superposition

See works by Barrau, De Lorenzo, Haggard, Christodoulou, Vilensky Rovelli, Speziale, Vidotto...

See also related works by Bianchi, Smerlak, Perez, Gosh, Frodden, Gambini, Pullin...

Bonanno, Reuter hep-th/0002196

Alkofer, D'Odorico, Saueressig, Vidotto 1503.06472



■ NON-PERTURBATIVE EFFECT

Effective theory: quantum repulsion

Quantum effects piling outside the horizon

■ TIME DILATATION

Bounce time $\sim M \sim \text{ms}$ for M_{\odot}

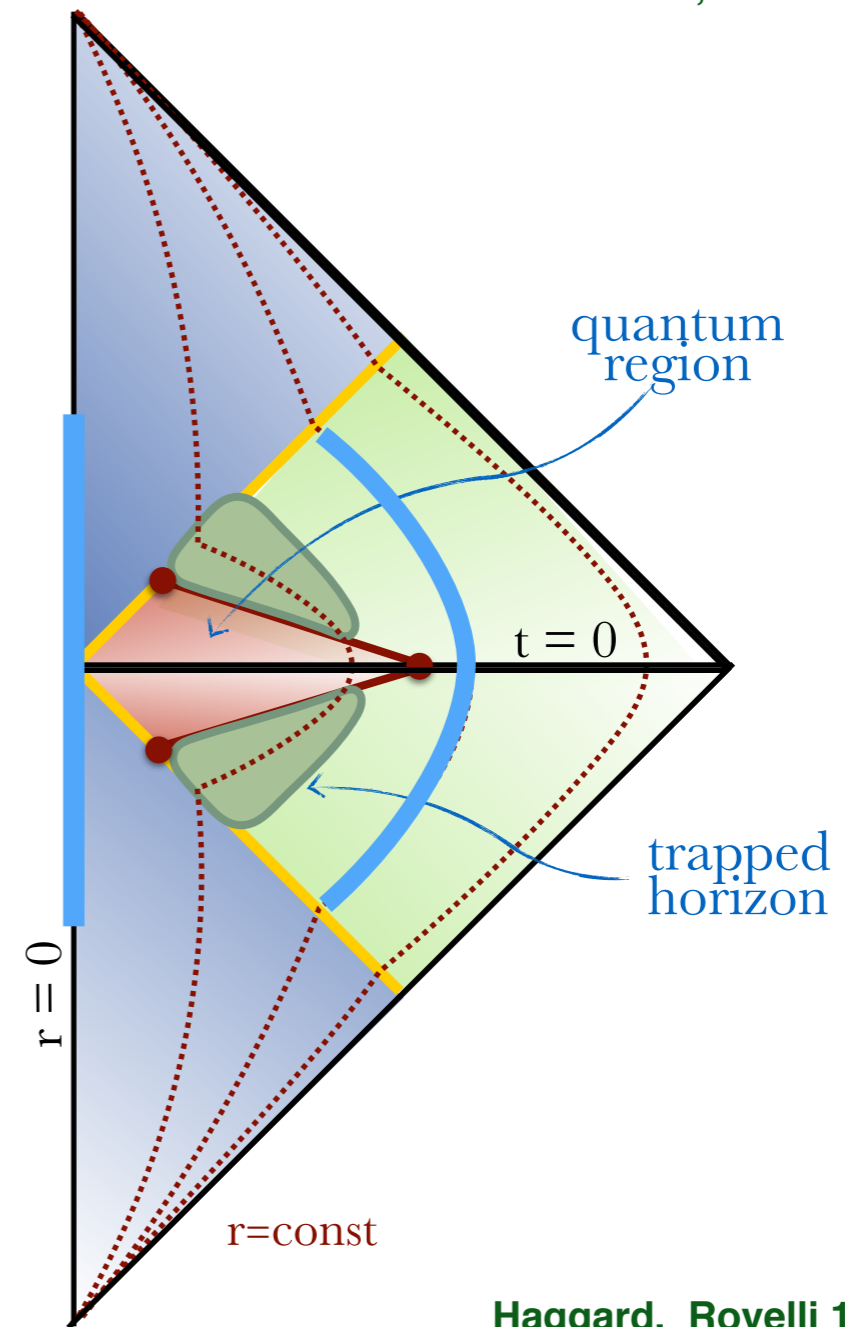
Asymptotic time $\sim M^2 \sim 10^9$ for M_{\odot}

■ LIFETIME $\sim M^2$

to be compared with the evaporation time $\sim M^3$

(no information paradox)

Vidotto, Rovelli 1401.6562



Haggard, Rovelli 1407.0989

■ Upper limit:

Vidotto, Rovelli 1401.6562

Firewall argument (Almheiri, Marolf, Polchinski, Sully): “something” unusual must happen before the Page time ($\sim 1/2$ evaporation time)

\Rightarrow the hole lifetime must be shorter or of the order of $\sim \mathbf{m^3}$

■ Lower limit:

Haggard, Rovelli 1407.0989

For something quantum to happen, semiclassical approximation must fail.

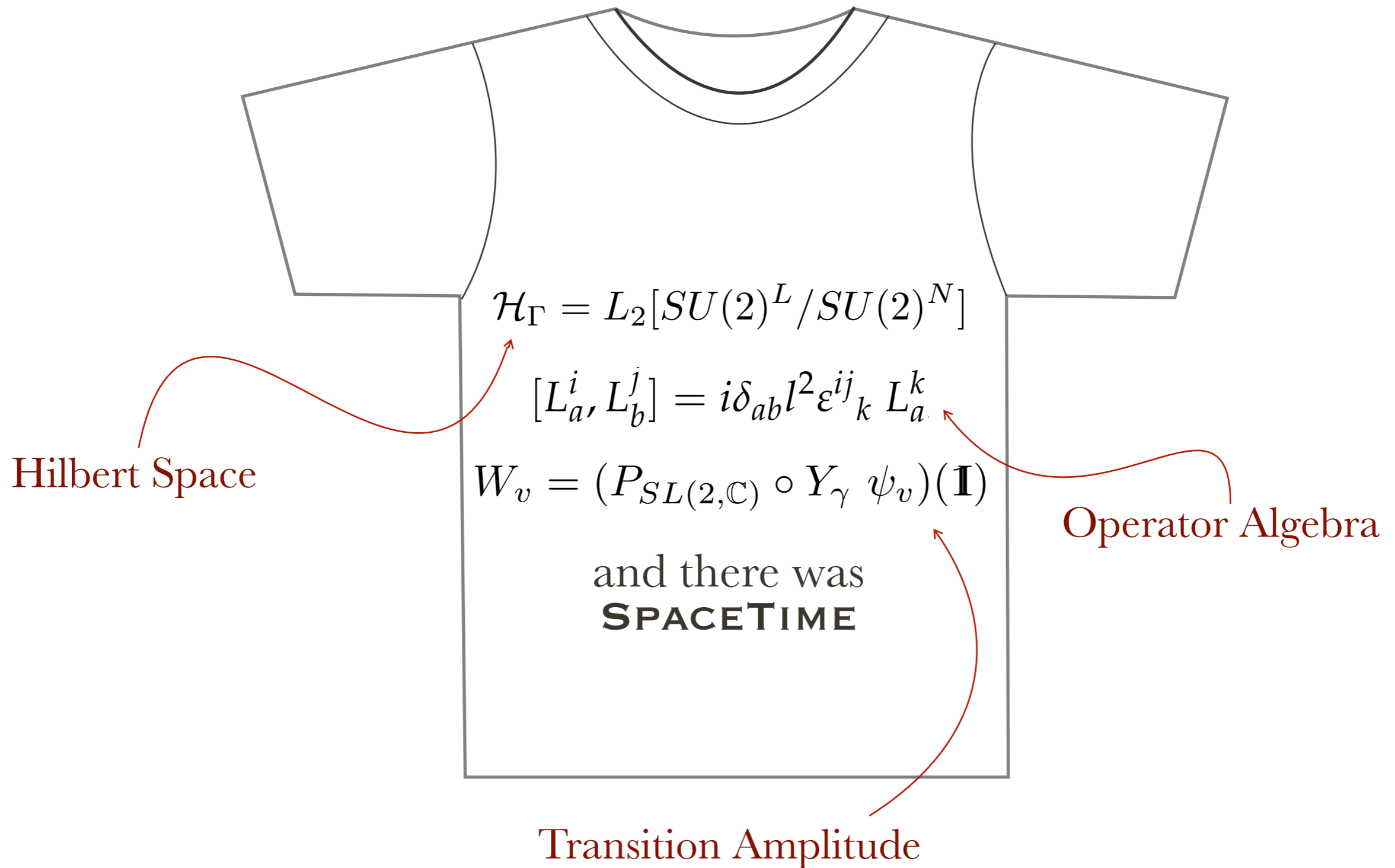
Typically in quantum gravity: high curvature **Curvature $\sim (\mathbf{L_P})^{-2}$**

Small effects can pile up: small probability per time unit gives a probable effect on a long time!

Typically in quantum tunneling: **Curvature \times (time) $\sim (\mathbf{L_P})^{-1}$**

$$\frac{1}{m^2} T_b \sim 1$$

\Rightarrow the hole lifetime must be longer or of the order of $\sim \mathbf{m^2}$

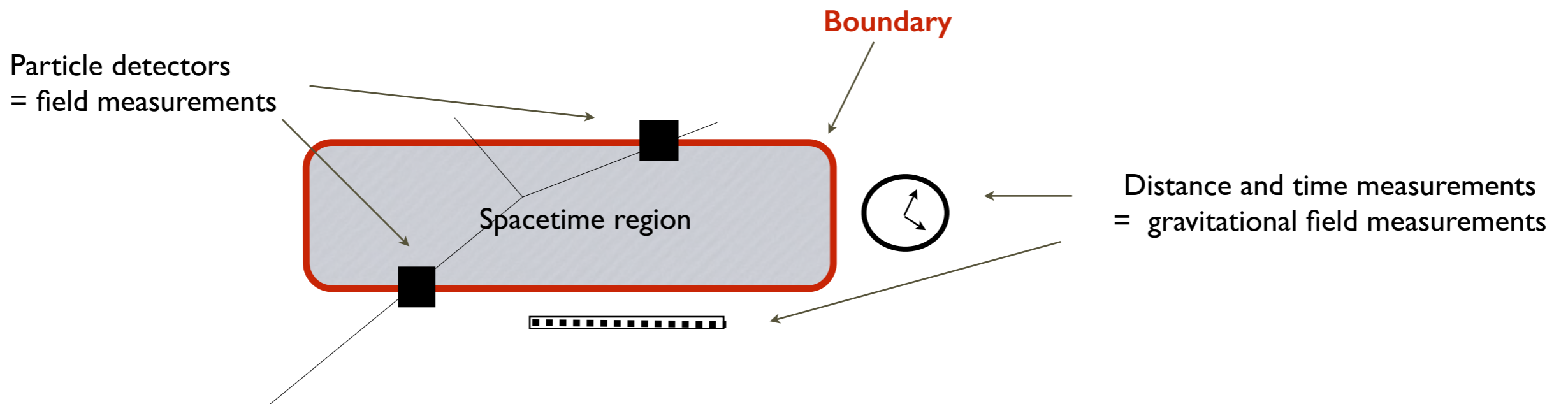


A PROCESS AND ITS AMPLITUDE

Boundary state $\Psi = \psi_{in} \otimes \psi_{out}$

Amplitude $A = W(\Psi)$

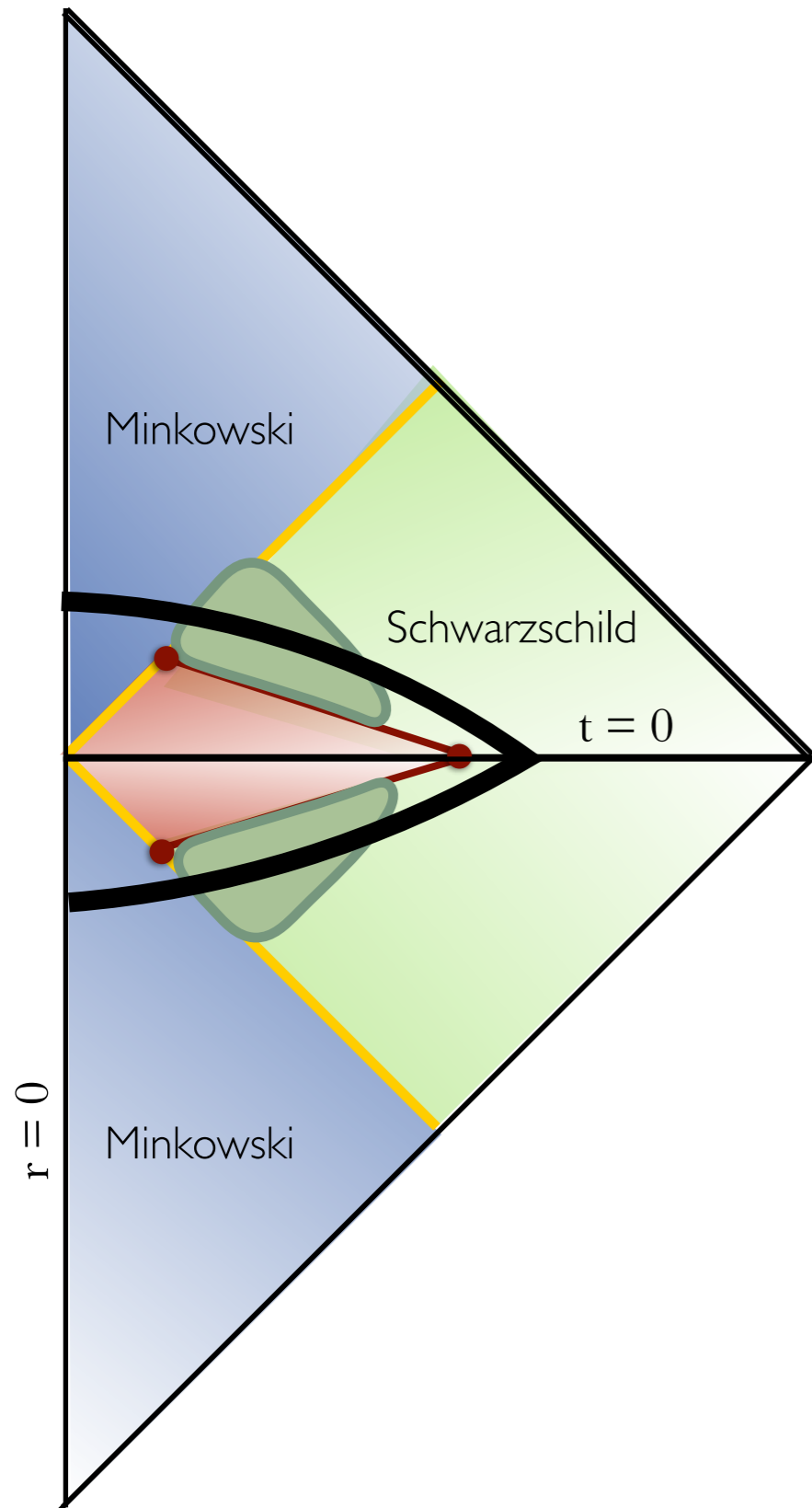
Quantum system
=
Spacetime region



In GR, distance and time measurements are field measurements like any other one:
they are part of the **boundary data** of the problem

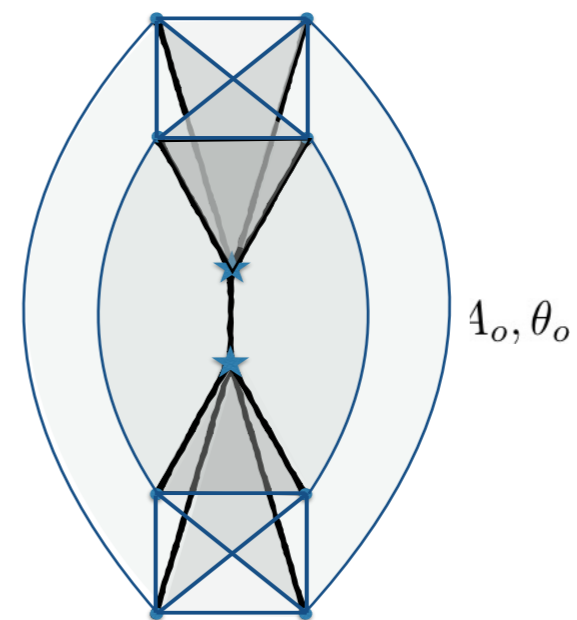
Boundary values of the gravitational field = geometry of box surface = distance and time separation of measurements

BOUNDARY STATE



- Boundary: $B_3 \cup B_3$ (joined on a S_2)
- Each B_3 can be triangulated by 4 isosceles tetrahedra
- The bulk can be approximated to first order by two 4-simplices joined by a tetrahedron

$$\Sigma = B_- \cup B_+$$



$$W(m, T) = \sum_{\{j_\ell\}} w(m, T, j_\ell) \sum_{\{J_n\}, \{K_n\}, \{l_\ell\}} \left(\bigotimes_n N_{\{j_n\}}^{J_n}(\{\nu_n\}, \{\alpha_n\}) f_{\{j_n\}\{l_n\}}^{J_n, K_n} \right) \left(\bigotimes_n i^{K_n, \{l_n\}} \right)_\Gamma$$

$$w(m, T, j_\ell) = c(m) \prod_\ell d_{j_\ell} e^{-\frac{1}{2\eta_\ell} \left(j_\ell - \frac{(2\eta_\ell^2 - 1)}{2} \right)^2} e^{i\gamma\zeta_\ell j_\ell}, \quad \eta_\ell^2 \sim m^2$$

$$T \sim m^2$$

$$f_{\{j_n\}\{l_n\}}^{K_n, J_n} \equiv d_{J_n} i^{J_n, \{j_n\}}_{\{\vec{p}_n\}} \left(\int dr_n \frac{\sinh^2 r_n}{4\pi} \bigotimes_{\ell \in n} \overrightarrow{d_{j_\ell l_\ell p_\ell}(r_n)} \right) i^{K_n, \{l_n\}}_{\{\vec{p}_n\}} d_{K_n}$$

$$\int_0^{\tau(m)} P(m, T) dT = 1 - \frac{1}{e}$$

- exploding now: $m(t)|_{t=t_H}$ $R = \frac{2Gm}{c^2}$

- **LOW ENERGY:** size of the source \approx wavelength $\lambda_{predicted}$
- **HIGH ENERGY:** energy of the particle liberated $\approx Tev$

- fast process (few milliseconds?)
- the source disappears with the burst
- very compact object: big flux $E = mc^2$

EXPONENTIAL DECAY: m^2 IS FAVORITE

■ exploding now: $m = \sqrt{\frac{t_H}{4k}} \sim 1.2 \times 10^{23} \text{ kg}$ $R = \frac{2Gm}{c^2} \sim .02 \text{ cm}$

- **LOW ENERGY:** size of the source \approx wavelength $\lambda_{predicted} \gtrsim .05 \text{ cm}$
- **HIGH ENERGY:** energy of the particle liberated $\approx Tev$

■ fast process (few milliseconds?)

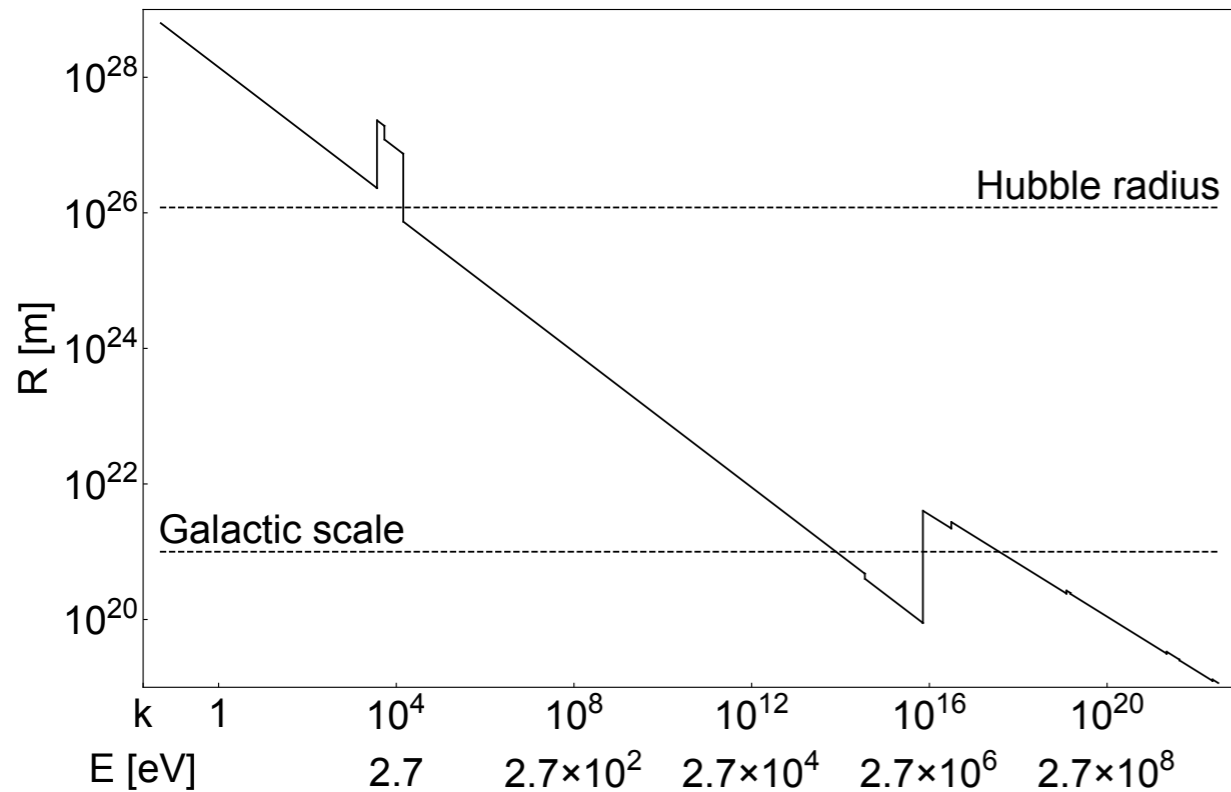
■ the source disappears with the burst

■ very compact object: big flux $E = mc^2 \sim 1.7 \times 10^{47} \text{ erg}$

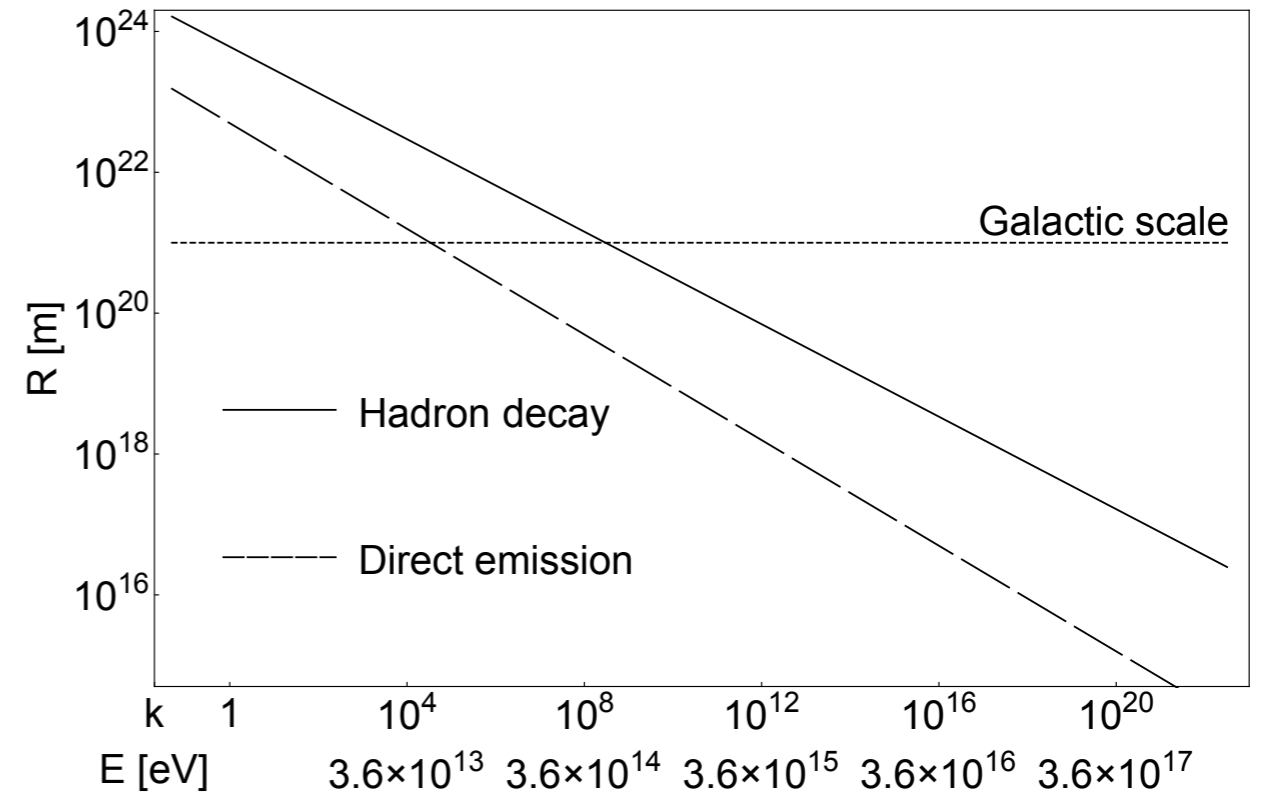
Barrau, Rovelli, Vidotto 1409.4031

■ shorter lifetime — smaller wavelength

Low energy channel



High energy channel

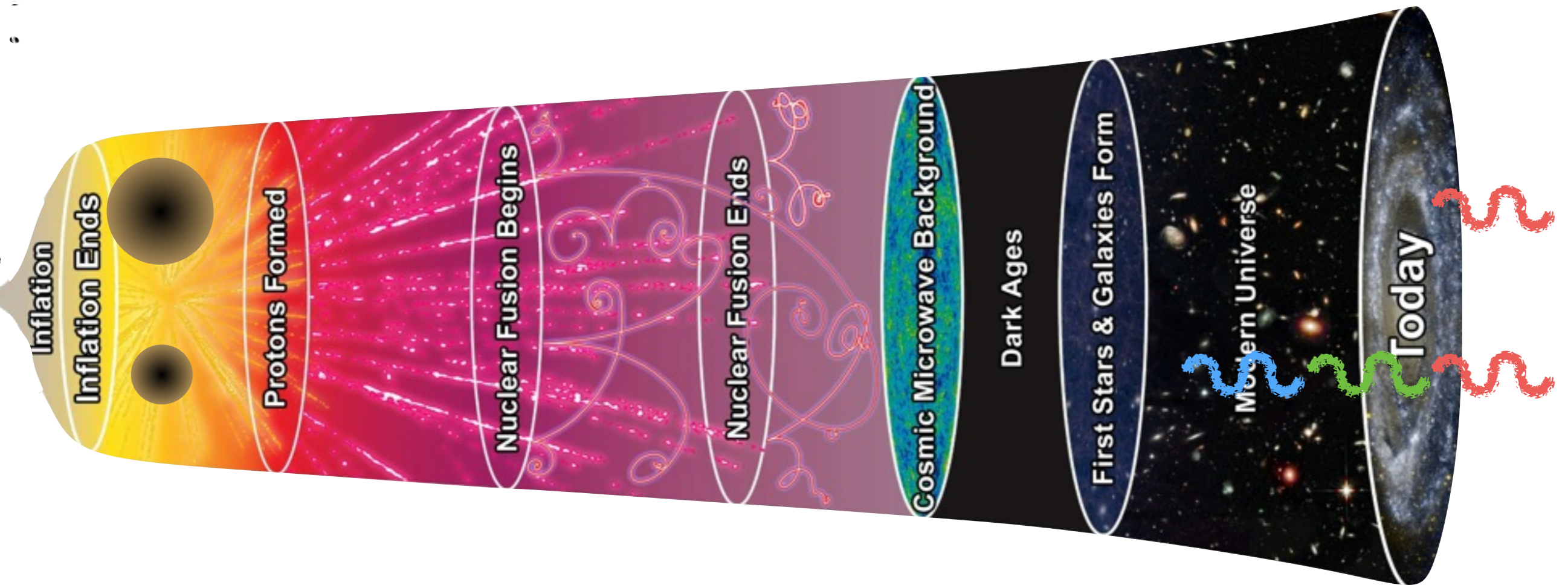


- detection of arbitrarily far signals
- better single-event detection

- PBH: mass - temperature relation
- different scaling

THE SMOKING GUN: DISTANCE/ENERGY RELATION

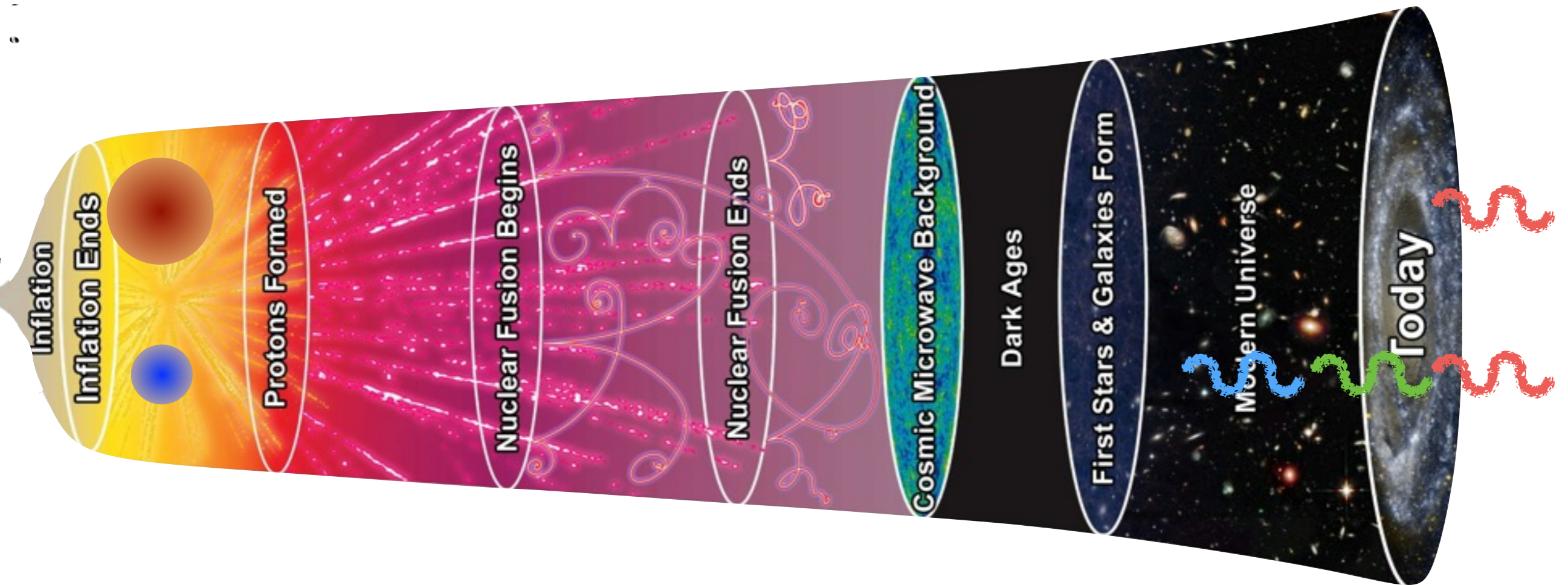
Low energy channel



- distant signals originated in younger and smaller sources

THE SMOKING GUN: DISTANCE/ENERGY RELATION

High energy channel



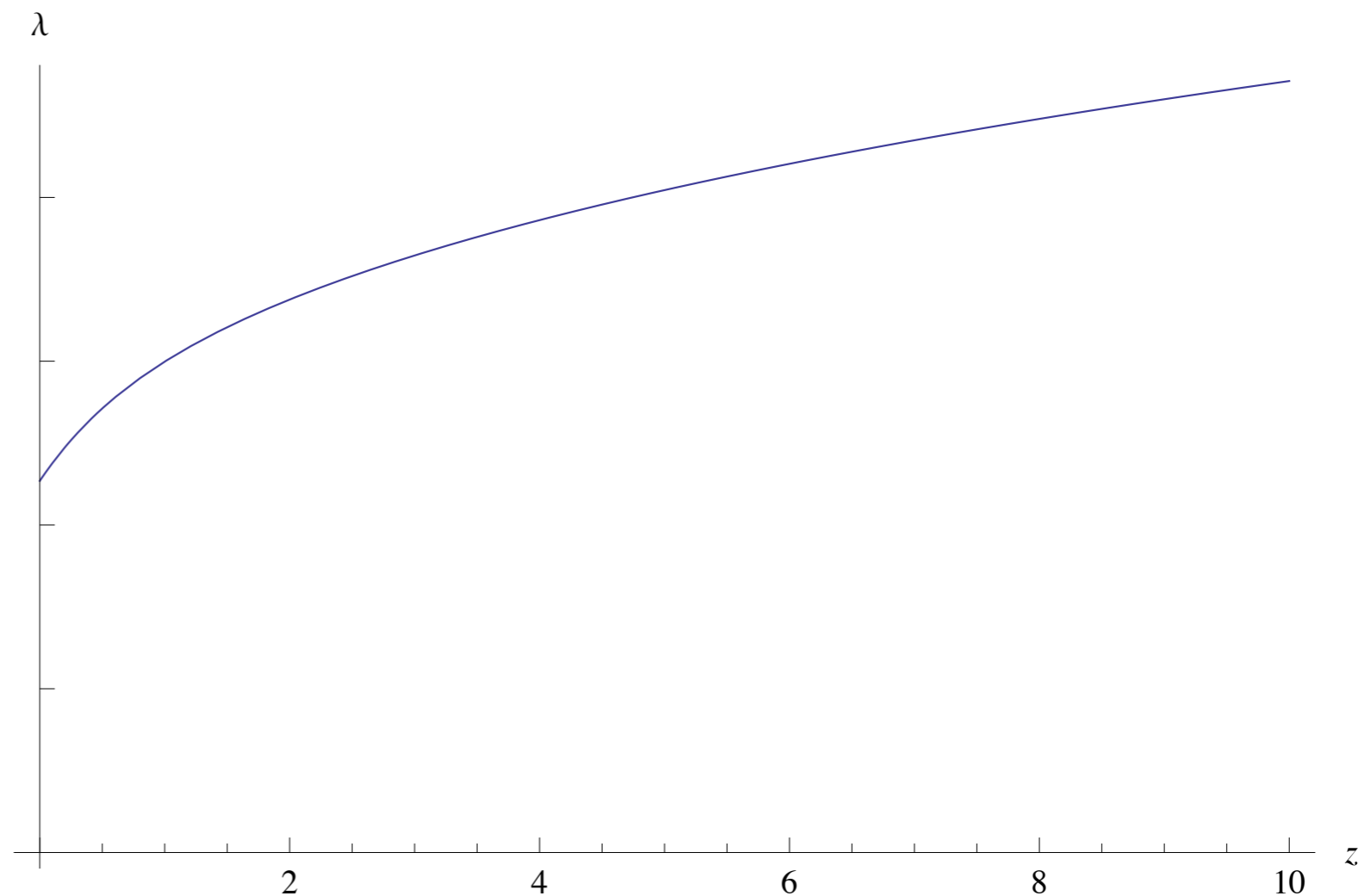
$$M \sim M_H \sim t.$$

$$t \sim 0.3g_*^{-\frac{1}{2}} T^{-2}$$

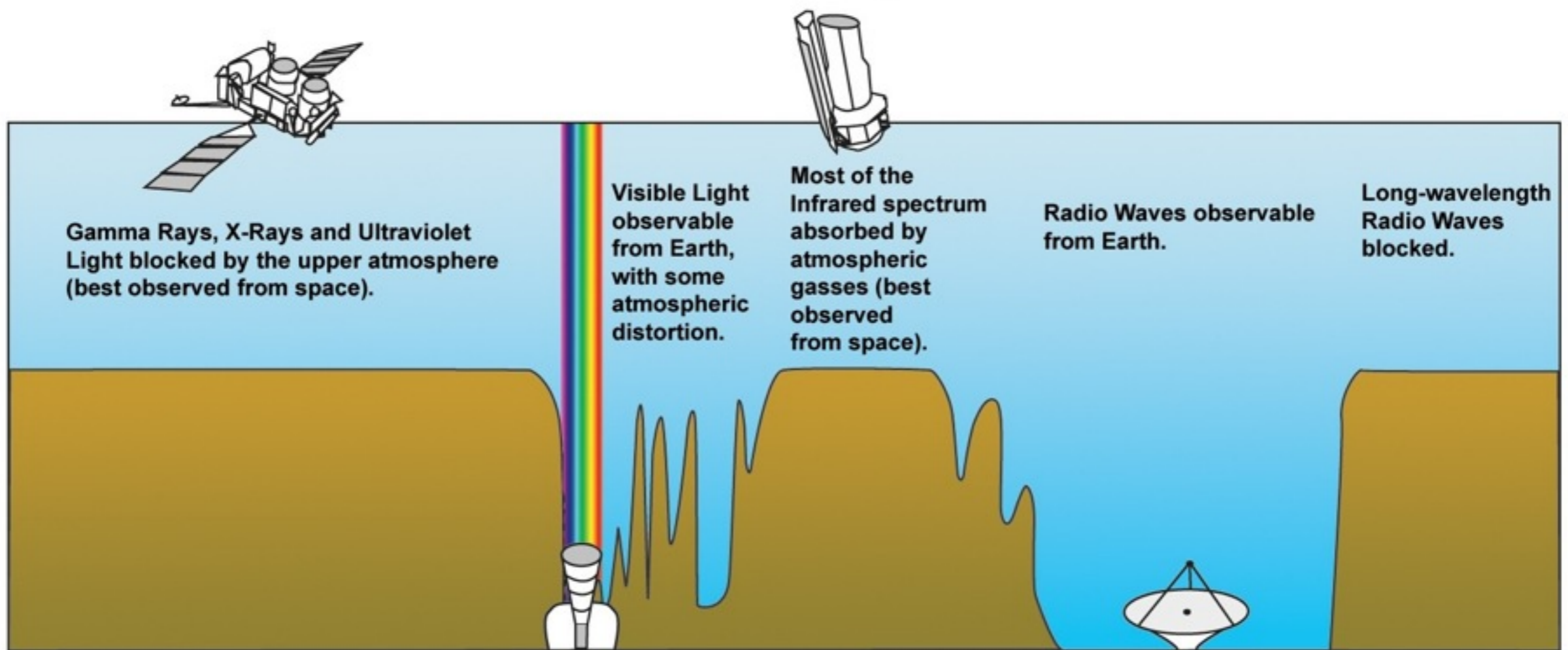
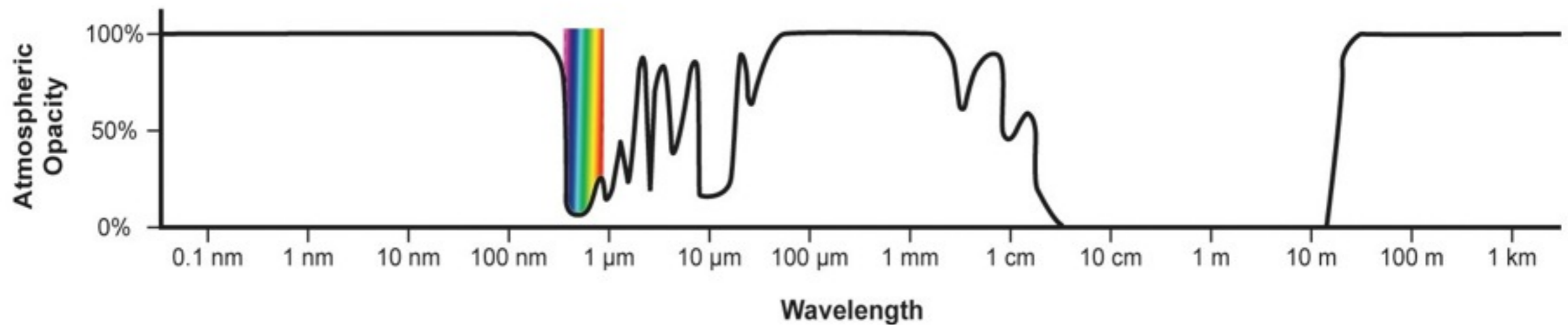
THE SMOKING GUN: DISTANCE/ENERGY RELATION

$$\lambda_{obs}^{other} = (1+z)\lambda_{emitted}^{other} \quad \longrightarrow \quad \lambda_{obs} \sim \frac{2Gm}{c^2}(1+z) \sqrt{\frac{H_0^{-1}}{6k\Omega_\Lambda^{1/2}} \sinh^{-1} \left[\left(\frac{\Omega_\Lambda}{\Omega_M} \right)^{1/2} (z+1)^{-3/2} \right]}$$

- distance $\propto 1/\text{wave length}$
- taking into account the **redshift** the resulting function is very slowly varying

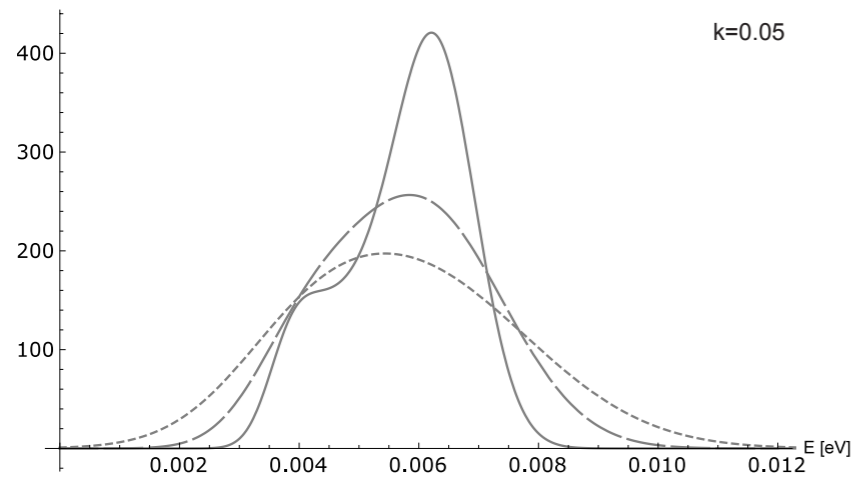


DETECTION ON EARTH?

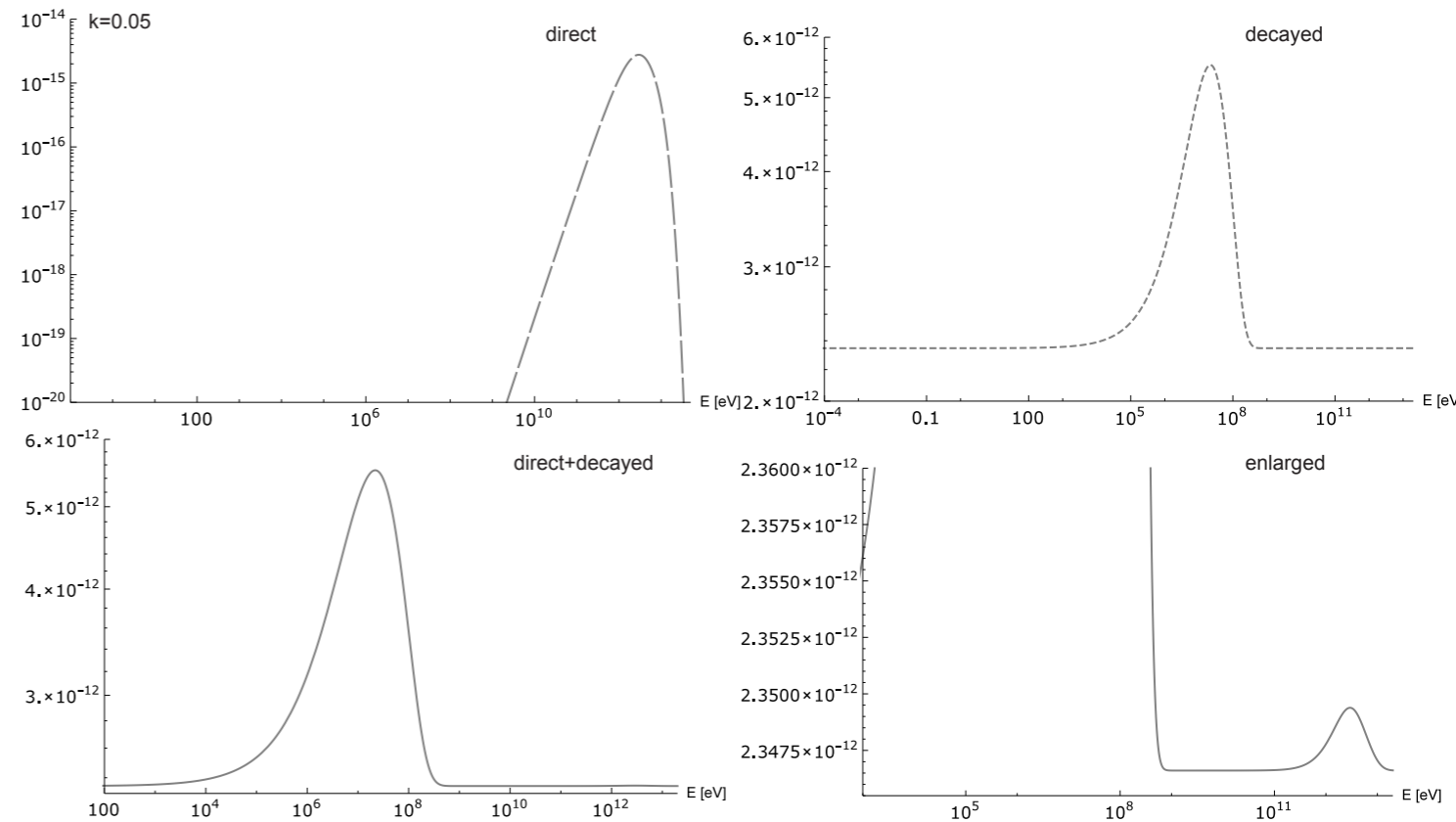


$$\tau \sim m^2$$

Low energy channel



High energy channel

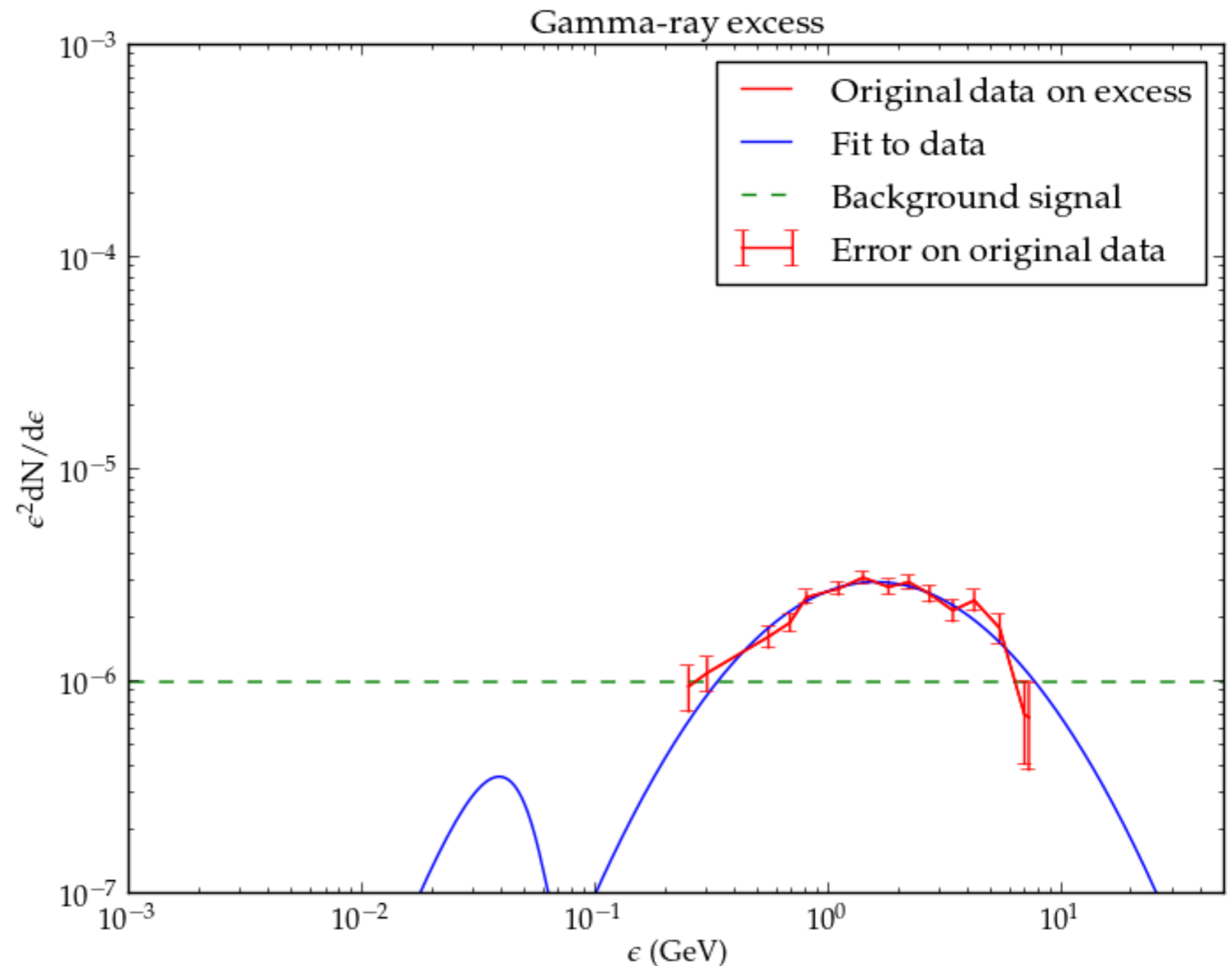


$$\frac{dN_{mes}}{dE dt dS} = \int \Phi_{ind}((1+z)E, R) \cdot n(R) \cdot Acc \cdot Abs(E, R) dR$$

- characteristic shape: distorted black body
- depends on how much DM are PBL

Low energy channel

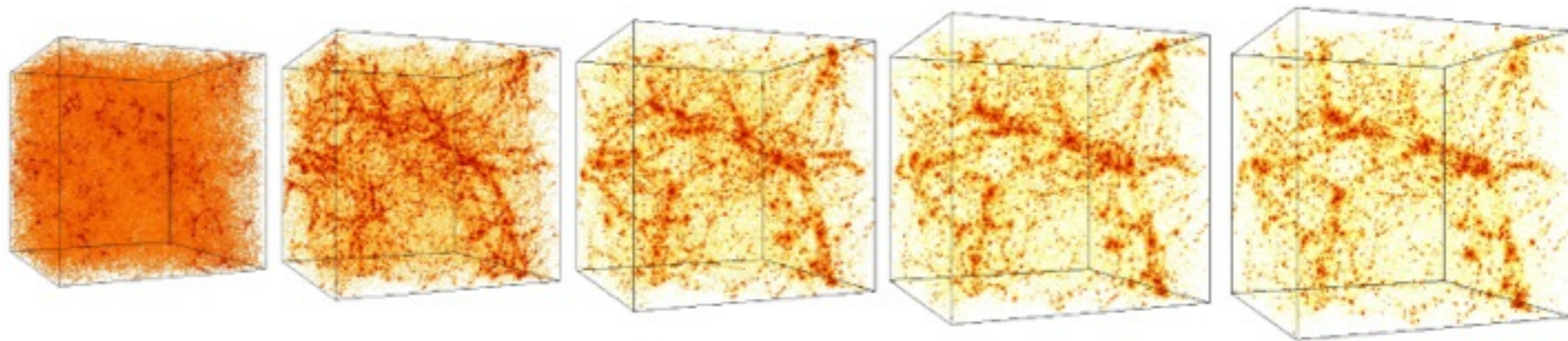
- Consider the longest possible lifetime of a quantum black hole.



- Number of secondary gamma-rays is higher than the number of primary gamma-rays, but their spectral energy density is much lower.

■ Structure formation

Raccanelli, Chluba, Cholis, Vidotto WIP

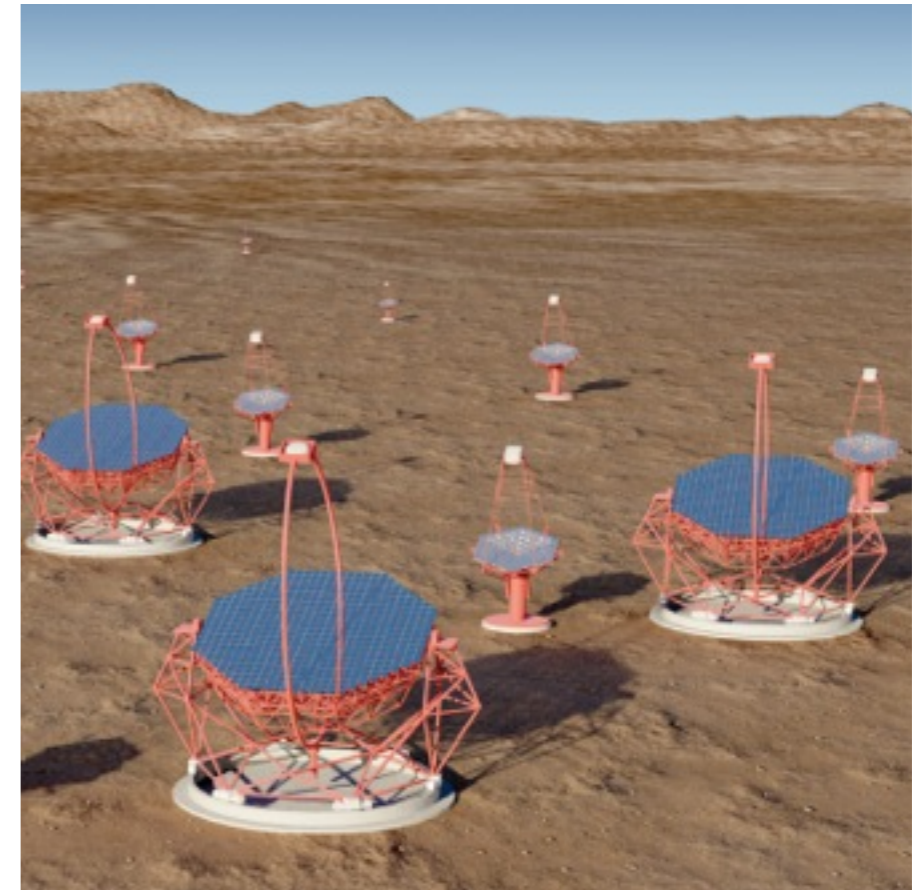


■ First stars & Supermassive black holes

Bambi, Freese, Vidotto WIP

- Primordial black holes inside first-generation stars can provide the seeds for supermassive black holes.

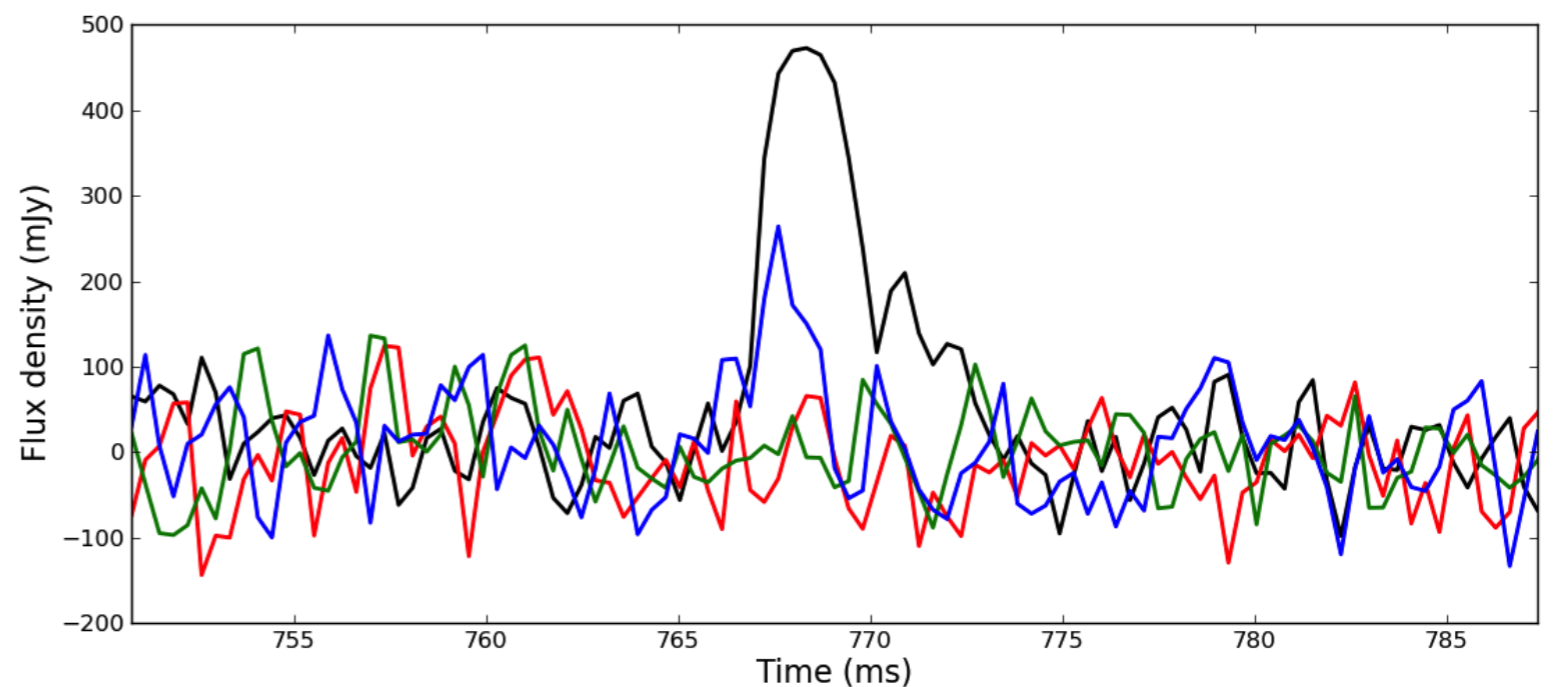
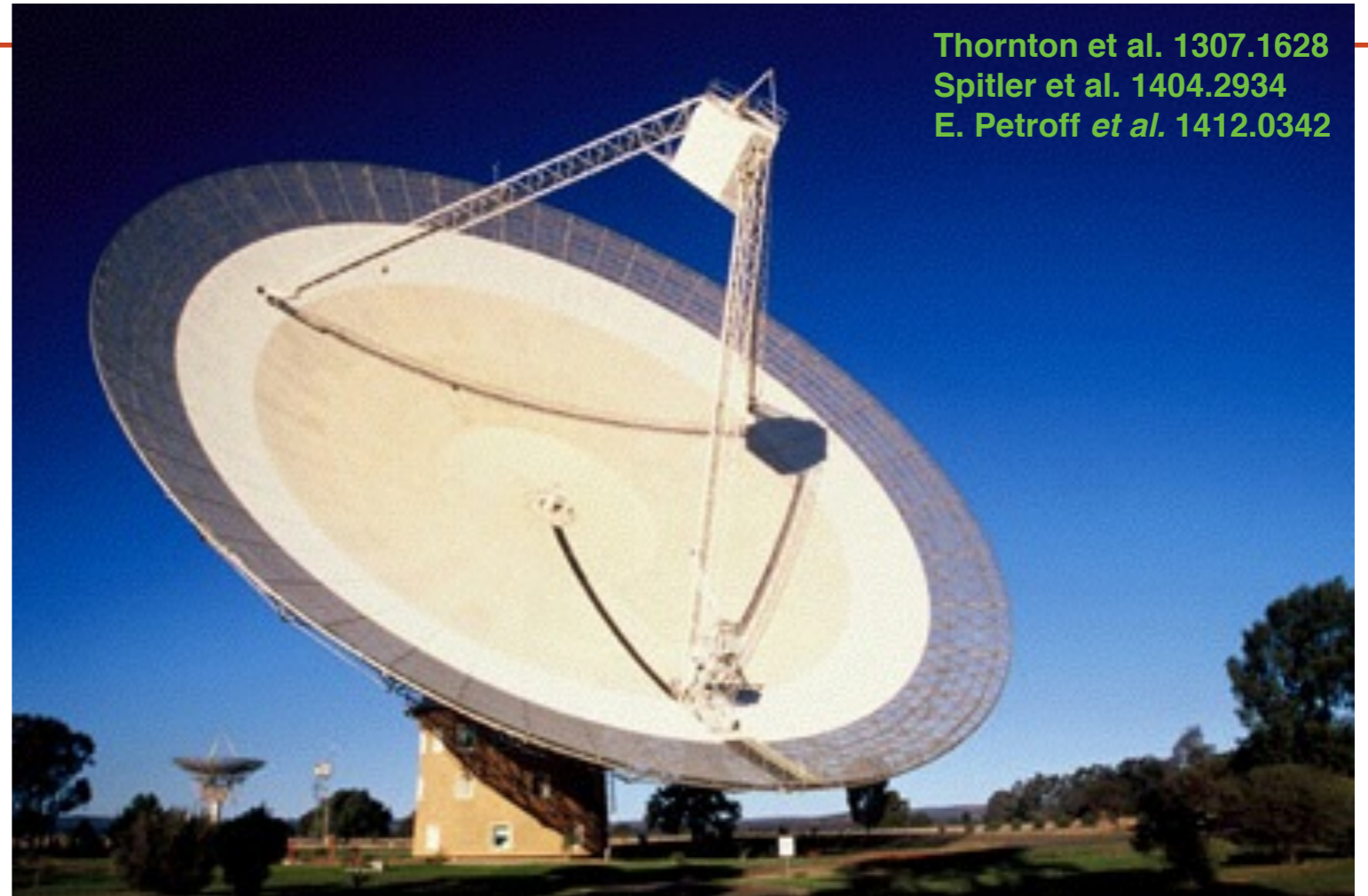
- the white hole should eject particles at the same temperature as the particles that fell in the black hole
- limited horizon due to absorption
~ 100 million light-years / $z=0.01$
- Short Gamma Ray Burst ?



- telescopes spanning large surfaces needed (CTA?)

Unknown source!

- Short
 - Observed width \approx milliseconds
- No Long GRB associated
 - No Afterglow
- Punctual
 - No repetition
- Enormous flux density
 - Energy $\approx 10^{38}$ erg
- Likely Extragalactic
 - Dispersion Measure: $z \lesssim 0.5$
- 10^4 event/day
 - A pretty common object?
- Circular polarization
 - Intrinsic




- $\lambda \approx 20$ cm \longrightarrow ■ size of the source $\approx \lambda_{predicted} \gtrsim .02$ cm
- Short
 - Observed width \approx milliseconds \longrightarrow ■ fast process
- No Long GRB associated
 - No Afterglow \longrightarrow ■ Very short GRB ? gravitational waves ?
- Punctual
 - No repetition \longrightarrow ■ the source disappears with the burst
- Enormous flux density
 - Energy $\approx 10^{38}$ erg \longrightarrow ■ very compact object $\longrightarrow 10^{47}$ erg
- Likely Extragalactic
 - Dispersion Measure: $z \lesssim 0.5$ \longrightarrow ■ peculiar distance/energy relation
- 10^4 event/day
 - A pretty common object?
- Circular polarization
 - Intrinsic

Are they bouncing Black Holes?

LIST OF FAST RADIO BURSTS

name	date	RA	dec	DM	width	peak	notes
FRB 010724	2001/07/24	01h18'	-75°12'	375	4.6	30	"Lorimer Burst"
FRB 010621	2001/06/21	18h52'	-08°29'	746	7.8	0.4	
FRB 110220	2011/02/20	22h34'	-12°24'	944.38	5.6	1.3	
FRB 110627	2011/06/27	21h03'	-44°44'	723.0	<1.4	0.4	
FRB 110703	2011/07/03	23h30'	-02°52'	1103.6	<4.3	0.5	
FRB 120127	2012/01/27	23h15'	-18°25'	553.3	<1.1	0.5	
FRB 011025	2001/10/25	19h07'	-40°37'	790	9.4	0.3	
FRB 121002	2012/10/02	18h14'	-85°11'	1628.76	2.1; 3.7	0.35	double pulse 5.1 ms apart
FRB 121002	2012/10/02	18h14'	-85°11'	1629.18	<0.3	>2.3	
FRB 121102	2012/11/02	05h32'	33°05'	557	3.0	0.4	by Arecibo radio telescope
	2015	05h32'~	33°05'~	557~			10 repeat bursts: 6 bursts in 10 minutes, 3 bursts weeks apart.
FRB 131104	2013/11/04	06h44'	-51°17'	779.0	<0.64	1.12	'near' Carina Dwarf Spheroidal Galaxy
FRB 140514	2014/05/14	22h34'	-12°18'	562.7	2.8	0.47	21 ± 7 per cent (3σ) circular polarization
FRB 090625	2009/06/25	03h07'	-29°55'	899.6	<1.9	>2.2	
FRB 130626	2013/06/26	16h27'	-07°27'	952.4	<0.12	>1.5	
FRB 130628	2013/06/28	09h03'	+03°26'	469.88	<0.05	>1.2	
FRB 130729	2013/07/29	13h41'	-05°59'	861	<4	>3.5	
FRB 110523	2011/05/23	21h45'	-00°12'	623.30	1.73	0.6	700-900 MHz at Green Bank radio telescope, detection of both circular and linear polarization.
FRB 150418	2015/04/18	07h16'	-19° 00'	776.2	0.8	2.4	Detection of linear polarization. The origin of the burst is disputed.

1. **BLACK HOLE** can be singularity free
and they can tunnel into a white hole in a time $\sim m^2$
* complete calculation available in LQG
2. **PHENOMENOLOGY** depends on mass and lifetime
* new experimental window for quantum gravity
 - IR radio & TeV
 - direct detection & diffuse emission
 - peculiar energy distance relation
3. **PRIMORDIAL BLACK HOLES**
 - new features



what else
can change if black holes
explode this way?

Planck Stars

Planck stars
[Carlo Rovelli](#), [Francesca Vidotto](#)
Int. J. Mod. Phys. D23 (2014) 12, 1442026

Classical metric

Black hole fireworks: quantum-gravity effects outside the horizon spark black to white hole tunneling
[Hal Haggard](#), [Carlo Rovelli](#)
Phys. Rev. D.92.104020.

Improved Black Hole Fireworks: Asymmetric Black-Hole-to-White-Hole Tunneling Scenario
[Tommaso De Lorenzo](#), [Alejandro Perez](#)
arXiv:1512.04566

Phenomenology

Planck star phenomenology
[Aurelien Barrau](#), [Carlo Rovelli](#).
Phys. Lett. B739 (2014) 405

Phenomenology of bouncing black holes in quantum gravity: a closer look
[Aurélien Barrau](#), [Boris Bolliet](#), [Francesca Vidotto](#), [Celine Weimer](#)
JCAP 1602 (2016) no.02, 022

Fast Radio Bursts

Fast Radio Bursts and White Hole Signals
[Aurélien Barrau](#), [Carlo Rovelli](#), [Francesca Vidotto](#).
Phys. Rev. D90 (2014) 12, 127503

LQG lifetime calculation

Computing a Realistic Observable in Background-Free Quantum Gravity: Planck-Star Tunnelling-Time from Loop Gravity
[Marios Chistodoulou](#), [Carlo Rovelli](#), [Simone Speziale](#), [Ilya Vilensky](#).
ArXiv: 1605.05268